

THE VISCOUS STRUCTURE OF BAROCLINIC CRITICAL LAYERS IN STRATIFIED SHEAR FLOWS WITH BACKGROUND ROTATION

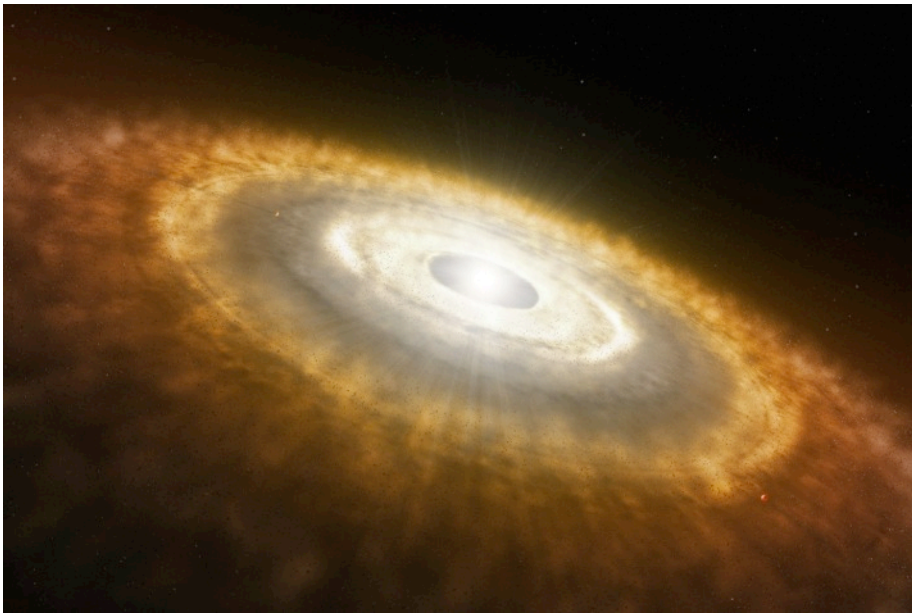
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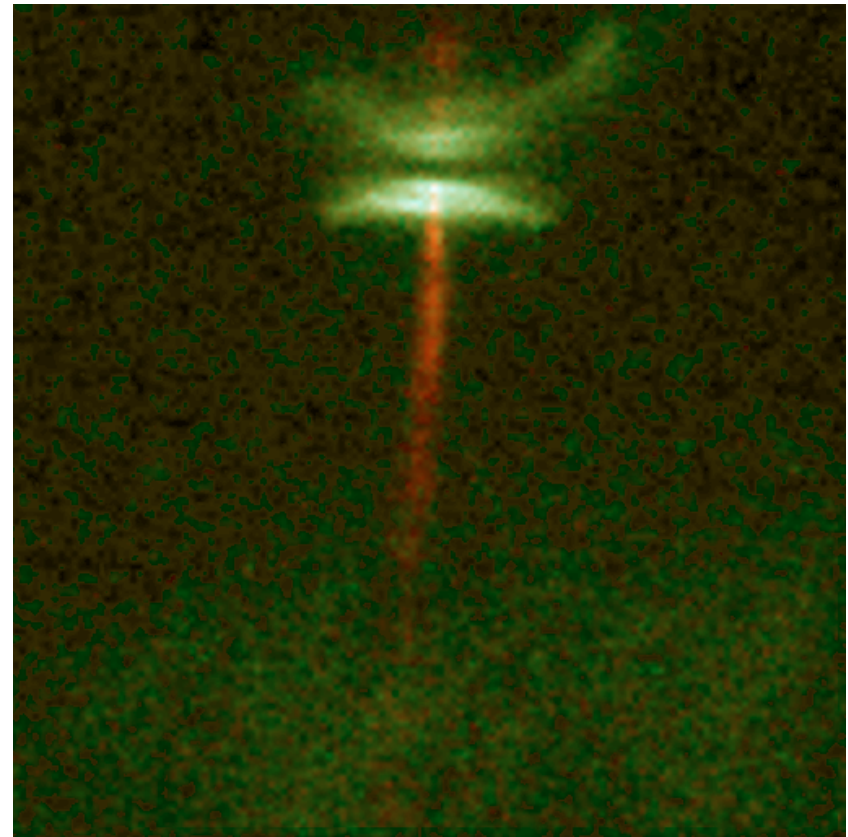
Emil Hopfinger Colloquium
May12, 2016



PROTO-PLANETARY DISKS (PPD): HOW DO STARS AND PLANETS FORM FROM PPD'S ?



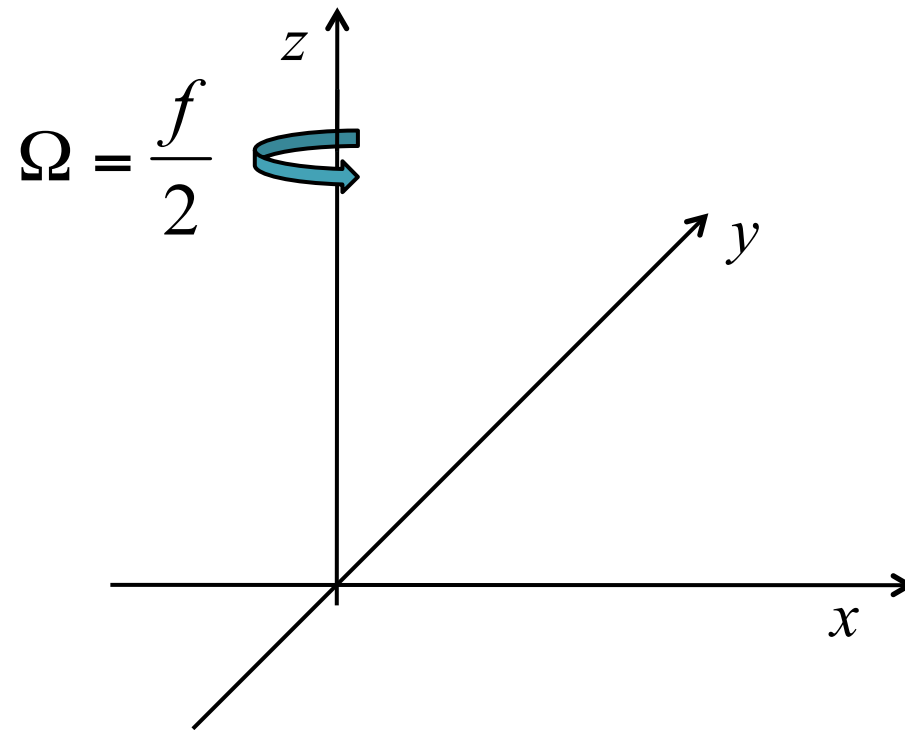
Artist's impression of “ baby star ”
surrounded by a proto-planetary disk



Proto-planetary disk of proto-star HH-30 in
Taurus constellation (ESA-NASA-Hubble)

ZOMBIE VORTICES

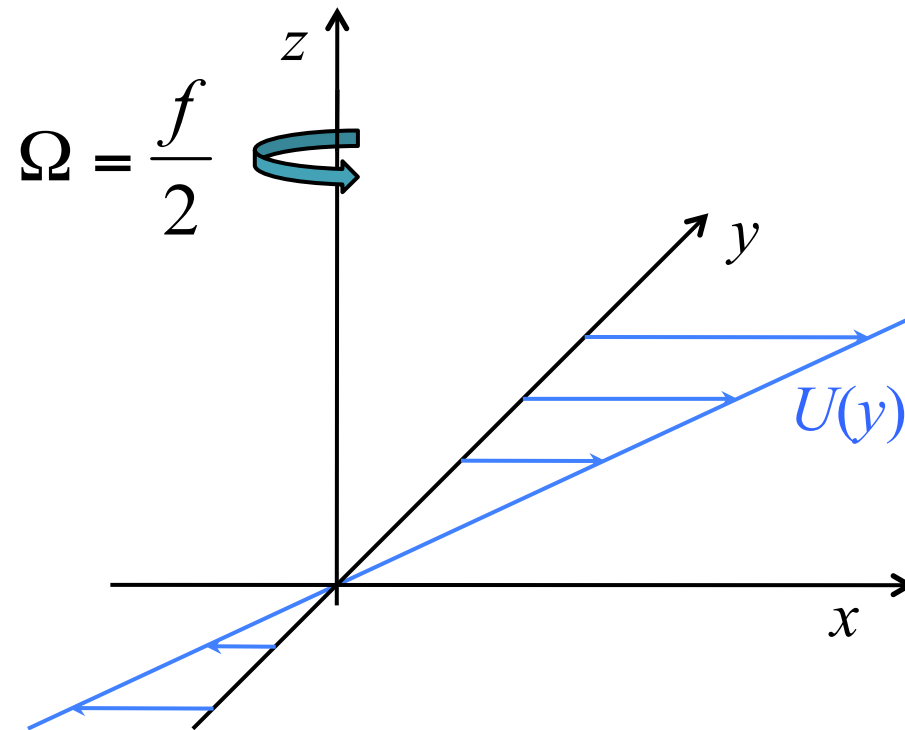
Idealized uniform shear flow schematic and baroclinic critical points



Marcus, Pei, Jiang & Hassanzadeh (2013)

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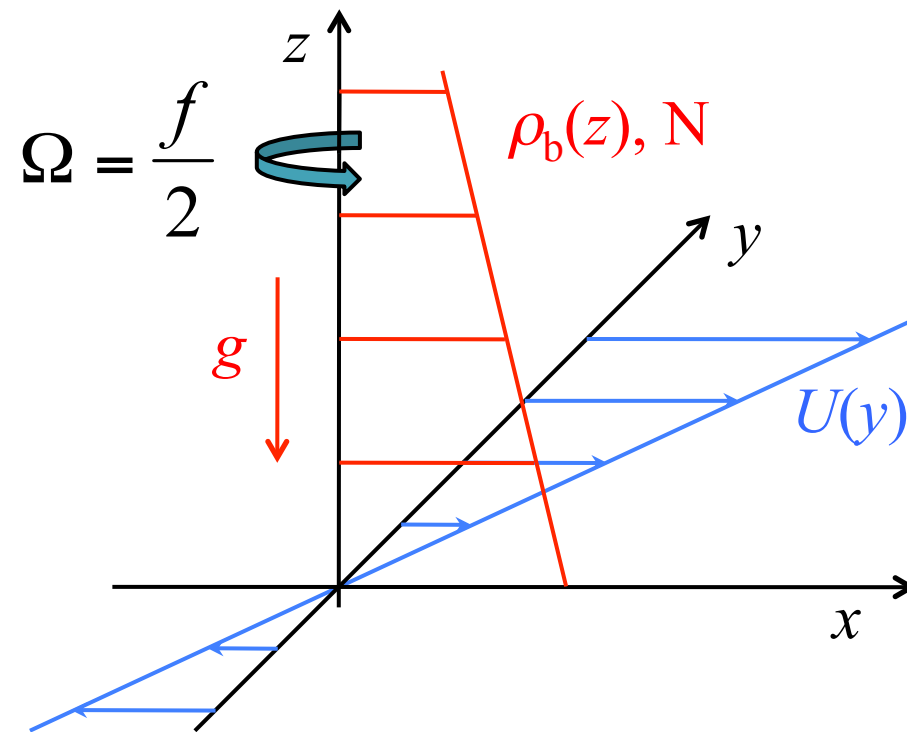
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ZOMBIE VORTICES

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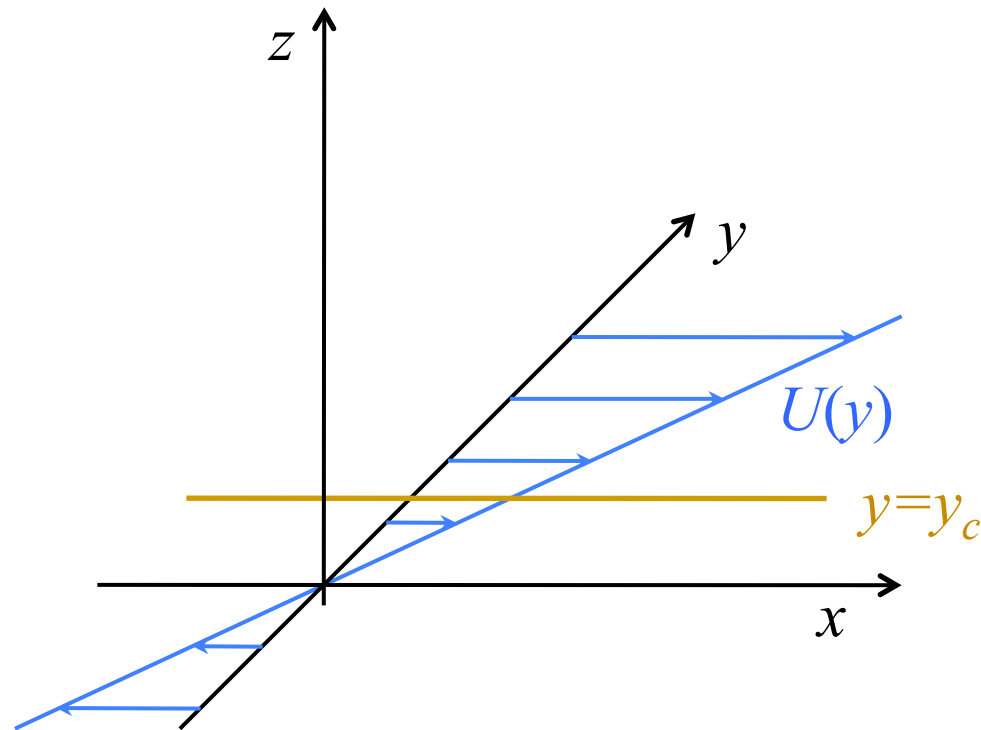


Marcus, Pei, Jiang & Hassanzadeh (2013)

ZOMBIE VORTICES

Idealized uniform shear flow schematic and baroclinic critical points

$$\omega - k_x U(y_c) = \pm N$$

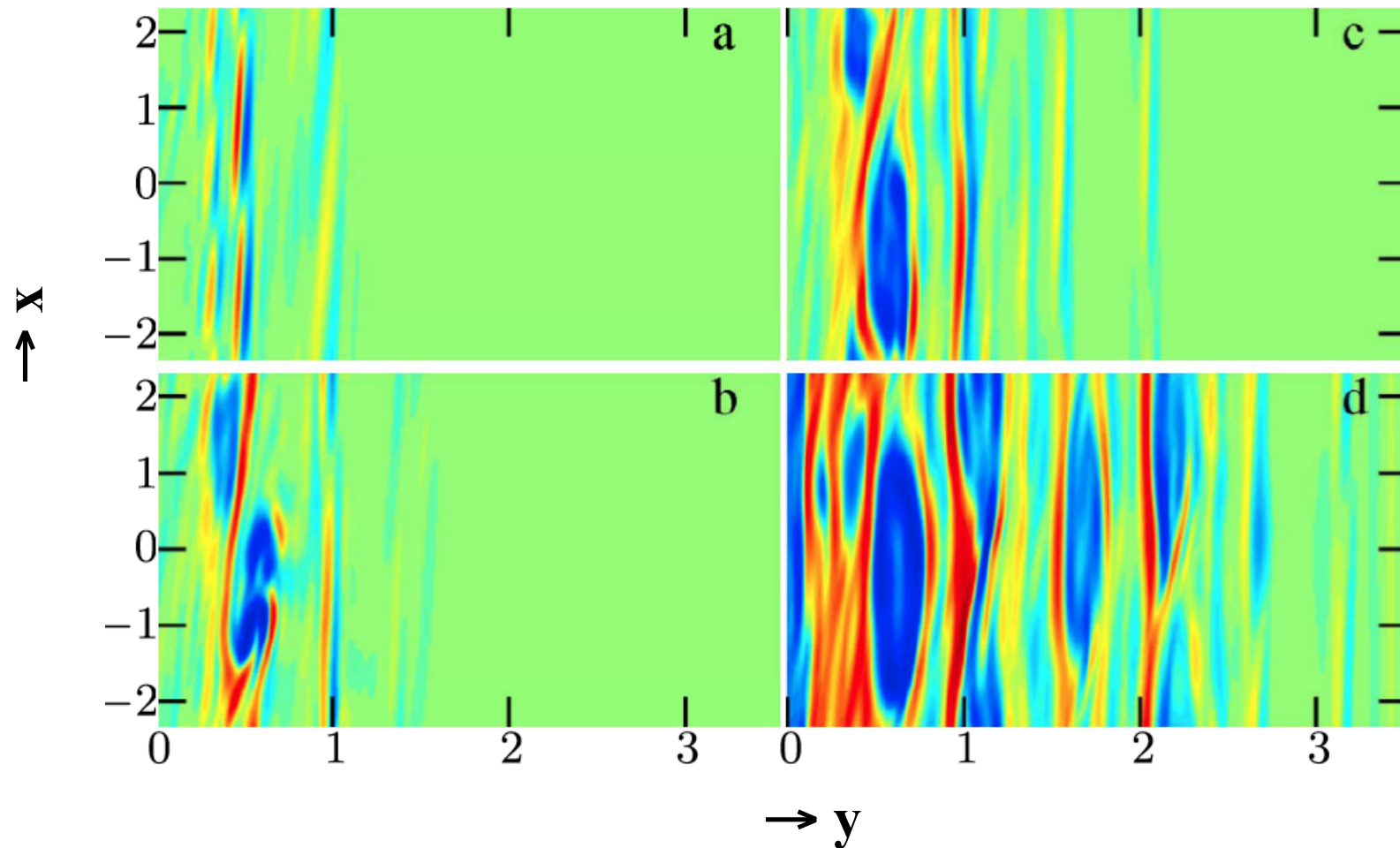


Marcus, Pei, Jiang & Hassanzadeh (2013)

ZOMBIE VORTICES

Inviscid numerical experiment:

vertical perturbation vorticity ω_z in horizontal x-y plane



Marcus, Pei, Jiang & Hassanzadeh (2013)

LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

Linear stability analysis of horizontal parallel shear flow with vertical linear stratification and background rotation

Horizontal

Froude number:

$$Fr_h = \frac{U_w}{Nh}$$

Rossby number:

$$Ro = \frac{U_w}{fh}$$

Reynolds number:

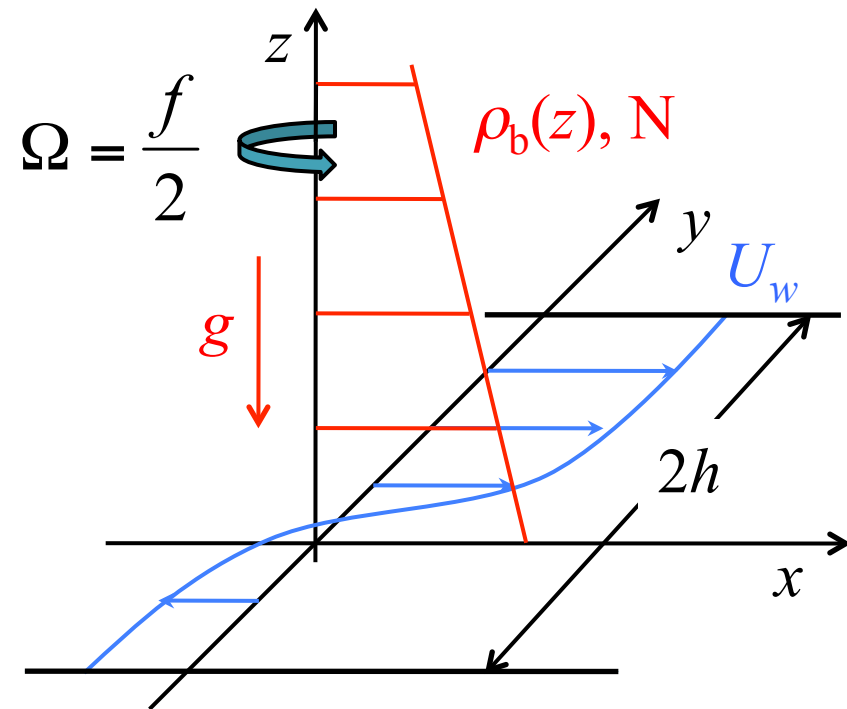
$$Re = \frac{U_w h}{\nu} \gg 1$$

Schmidt number:

$$Sc = \frac{\nu}{\kappa}$$

Normal mode:

$$q(x, y, z, t) = \Re \left\{ \hat{q}(y) e^{i(k_x x + k_z z - \omega t)} \right\}$$



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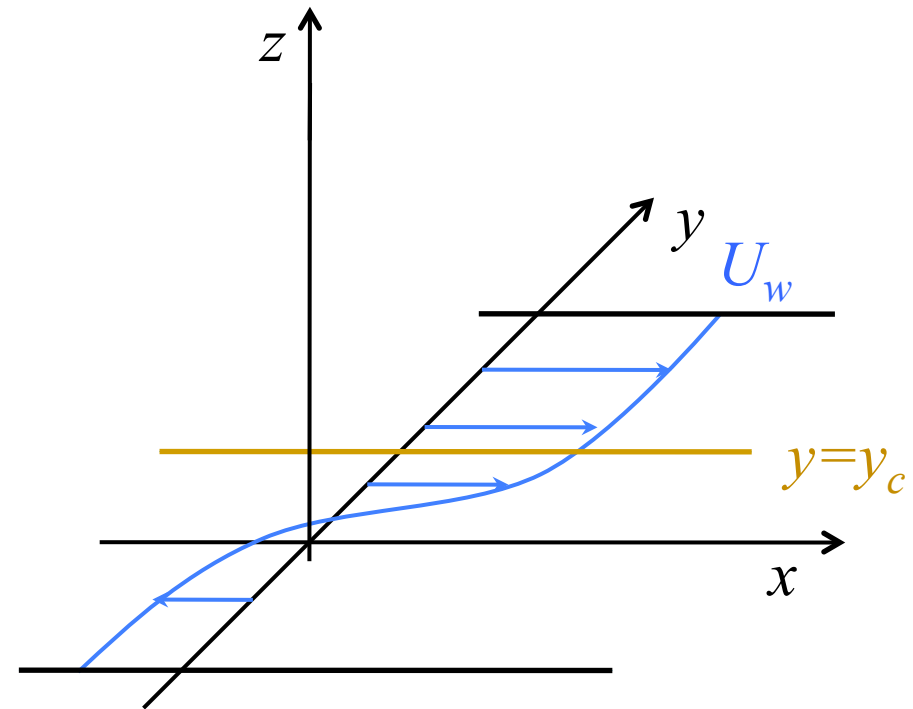
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LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

Inviscid “Rayleigh-like” outer region Nature of critical point singularity

$$\text{Baroclinic critical points: } U(y_c) - c = \pm \frac{1}{Fr_h k_x}$$

When $|y - y_c| \ll 1$:

$$\hat{u} = \frac{(U'(y_c) - \frac{1}{Ro}) a_0 \ln(y - y_c) - (U'(y_c) - \frac{1}{Ro}) (c_1 - a_0) - k_x \frac{1}{Fr_h} c_0}{\frac{1}{Fr_h^2} + \frac{1}{Ro} (U'(y_c) - \frac{1}{Ro})} + \dots,$$

$$\hat{v} = \frac{-i \frac{1}{Fr_h} a_0 \ln(y - y_c) + i \frac{1}{Fr_h} (c_1 - a_0) - i k_x \frac{1}{Ro} c_0}{\frac{1}{Fr_h^2} + \frac{1}{Ro} (U'(y_c) - \frac{1}{Ro})} + \dots,$$

$$\hat{w} = -\frac{k_z}{2k_x U'(y_c)} \frac{c_0}{y - y_c} + \dots,$$

$$\hat{\rho} = i Fr_h \frac{k_z}{2k_x U'(y_c)} \frac{c_0}{y - y_c} + \dots$$

LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

Inner viscous critical layer structure

“Smoothing” of critical point singularity by viscosity and density diffusion

“Zooming” inner cross-stream variable: $\eta = \frac{y - y_c}{\epsilon}, \quad \epsilon = i(k_x Re)^{-\frac{1}{3}}$

Vertical velocity:
$$\hat{w}^*(\eta) = \frac{\pi k_z}{k_x} \left(1 + \frac{1}{Sc}\right)^{-\frac{1}{3}} [2U'(y_c)]^{-\frac{2}{3}} \hat{p}^* Hi(\xi)$$

with
$$\xi = \text{sgn}[U'(y_c)] \left(\frac{2|U'(y_c)|}{1 + Sc^{-1}}\right)^{\frac{1}{3}} \eta$$

Scorer function $Hi(\xi)$:
$$\frac{d^2 Hi}{d\xi^2} - \xi Hi = \frac{1}{\pi}$$

LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

Inner viscous critical layer structure

“Smoothing” of critical point singularity by viscosity and density diffusion

Density:
$$\hat{\rho}^*(\eta) = -i\pi \frac{k_z}{k_x} Fr_h \left(1 + \frac{1}{Sc}\right)^{-\frac{1}{3}} [2U'(y_c)]^{-\frac{2}{3}} \hat{p}^* Hi(\xi)$$

Cross-stream velocity:
$$\hat{v}^*(\eta) = -i\pi \frac{k_z^2/k_x}{2U'(y_c)} \hat{p}^* \int_0^\xi Hi(t) dt + b^*$$

Similarly for $\hat{u}^*(\eta)$.

LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

Jump conditions across critical point between y_{c+} and y_{c-}

Streamwise velocity:

$$\hat{u}(y_{c+}) - \hat{u}(y_{c-}) = -i\pi s F r_h k_z^2 \hat{p}(y_c) \frac{U'(y_c) - 1/Ro}{2k_x U'(y_c)}$$

Cross-stream velocity:

$$\hat{v}(y_{c+}) - \hat{v}(y_{c-}) = -\frac{\pi s k_z^2 \hat{p}(y_c)}{2k_x U'(y_c)}$$

Reynolds stress:

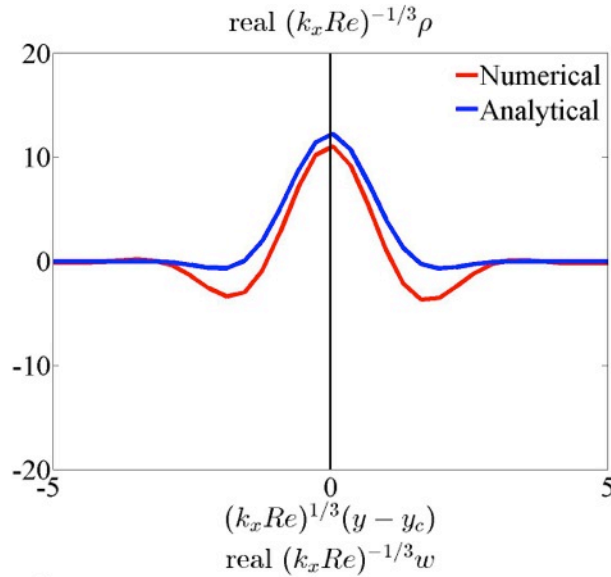
$$\tau_{xy}(y_{c+}) - \tau_{xy}(y_{c-}) = -\pi s F r_h k_z^2 \frac{|\hat{p}(y_c)|^2}{4U'(y_c)}$$

LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

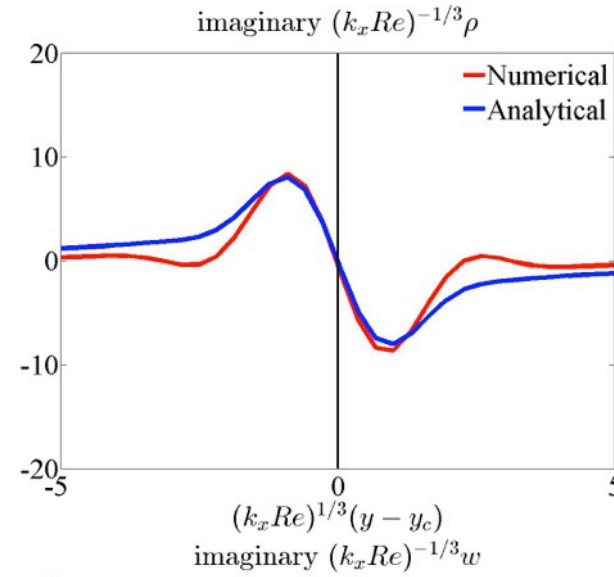
Inner viscous critical layer structure

Cross-stream distributions of density ρ^* and vertical velocity w^*

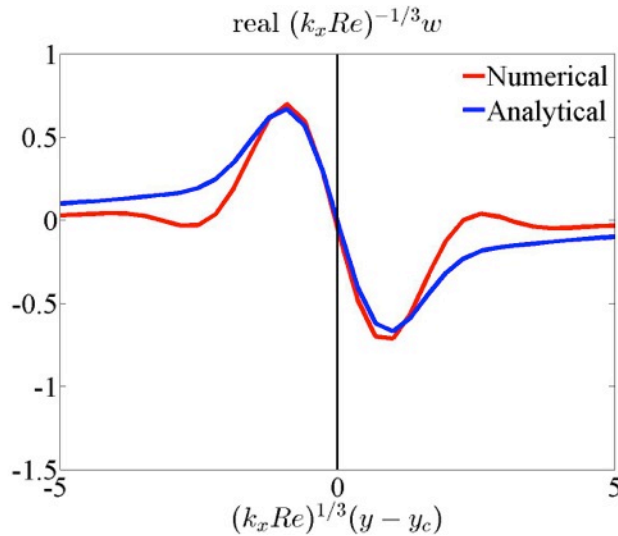
$$\frac{\rho_r^*}{\sqrt[3]{k_x Re}}$$



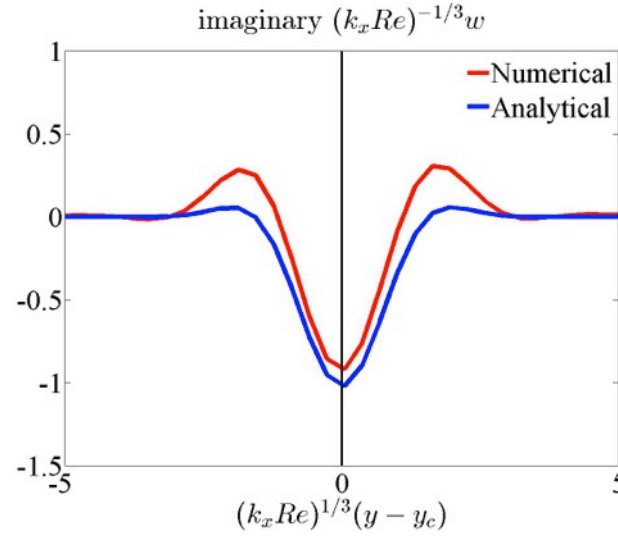
$$\frac{\rho_i^*}{\sqrt[3]{k_x Re}}$$



$$\frac{w_r^*}{\sqrt[3]{k_x Re}}$$



$$\frac{w_i^*}{\sqrt[3]{k_x Re}}$$

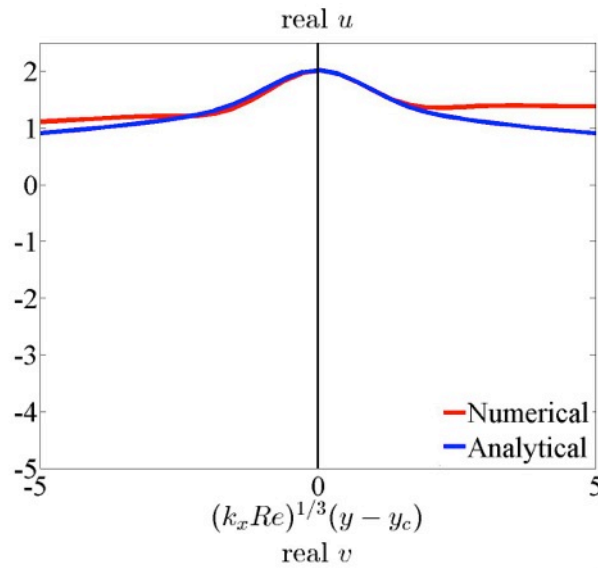


LARGE REYNOLDS NUMBER CRITICAL LAYER ANALYSIS

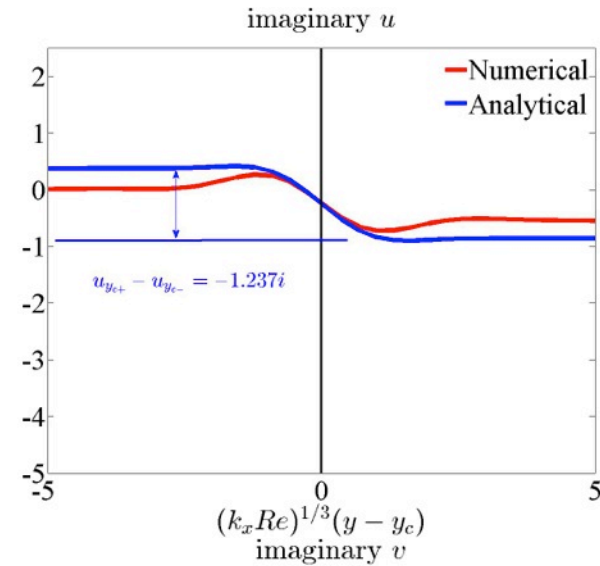
Inner viscous critical layer structure

Cross-stream distributions of streamwise velocity u^*
and cross-stream velocity v^*

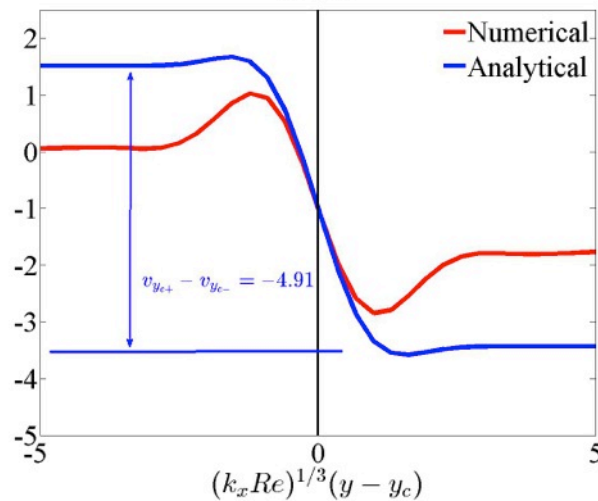
u_r



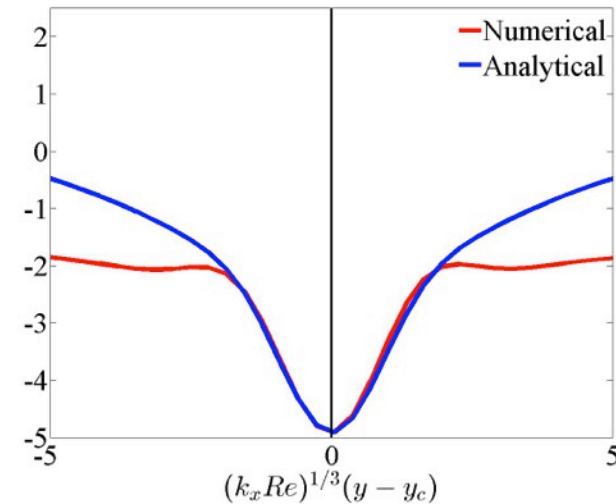
u_i



v_r



v_i



CONCLUSIONS AND ONGOING WORK

Horizontal shear flows in vertically stratified and rotating flows generate *baroclinic* critical points that are more singular than classical 2D *barotropic* critical points

The structure of *baroclinic critical layers* has been determined in the *linear and viscous régime*.

Strong jet-like and shear-layer-like vertical velocity profiles are produced which may lead to large-scale vortices (Boulanger et al. 2007; 2008) A Floquet instability analysis will tell whether these periodic flows give rise to such structures.

Ongoing investigation of the *nonlinear and viscous critical layer régime*