

Evidence of Görtler vortices in katabatic jet along a convexly curved slope.

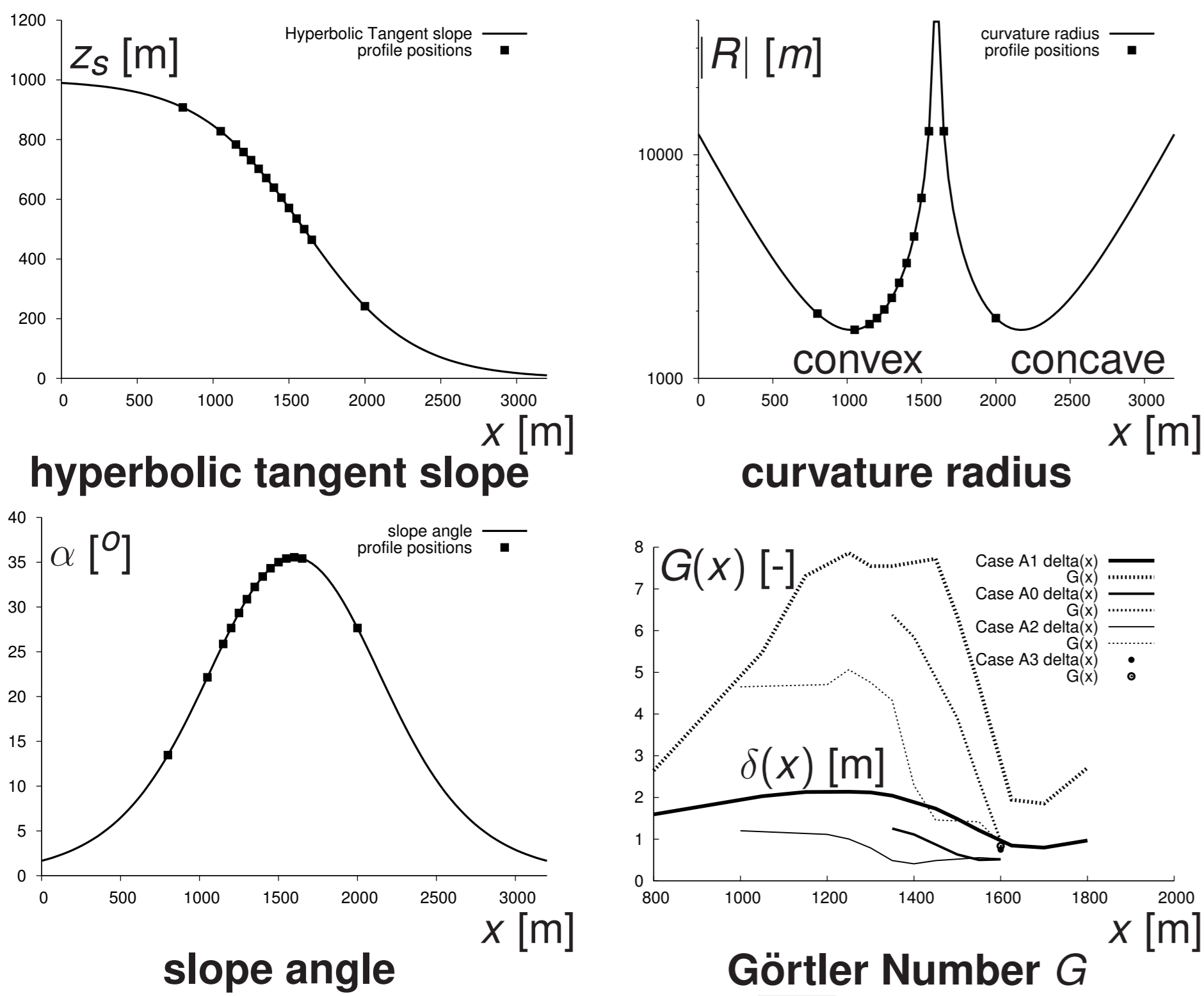
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Abstract

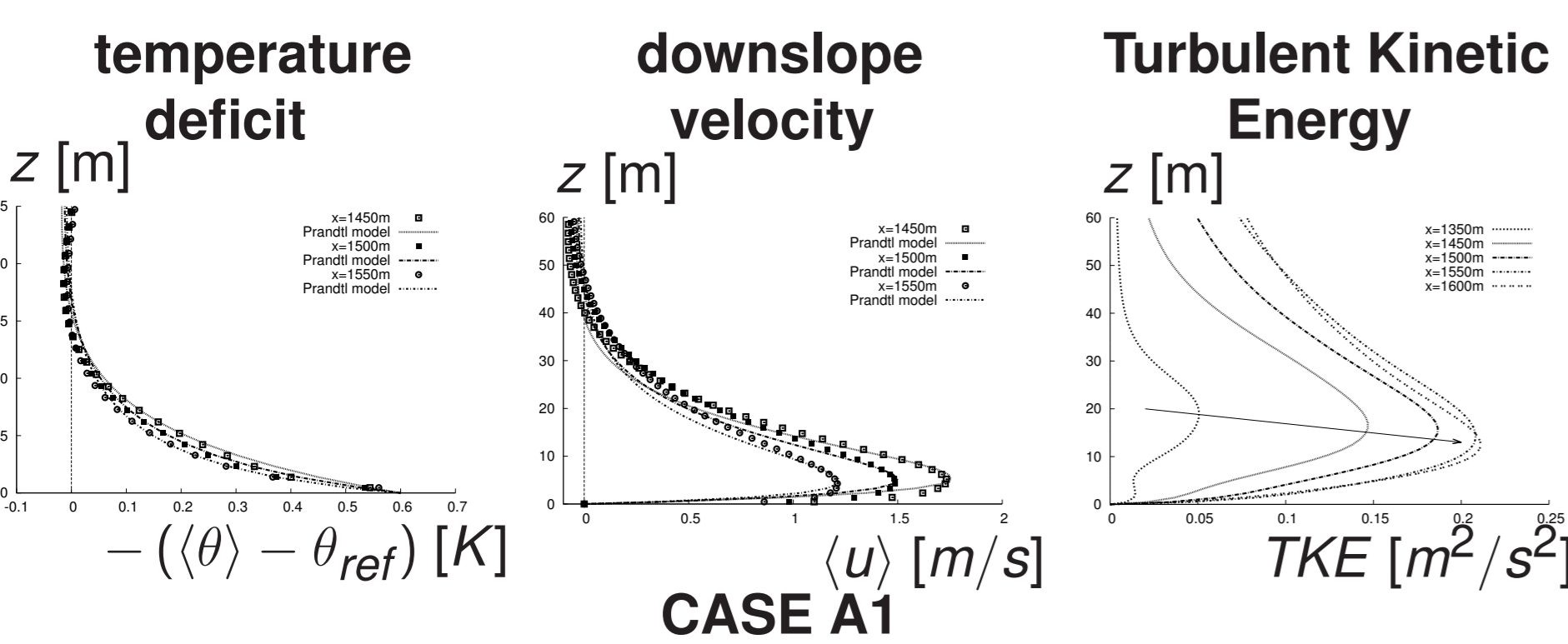
Large Eddy Simulation of katabatic flow along a convexly curved slope is performed. A special focus is given on the outer-layer shear of the katabatic jet. Both a statistical quantitative analysis and a qualitative description of vortical structures are used to describe the present turbulent flow. It is shown that Görtler vortices oriented in the streamwise downslope direction and with a vertical mushroom shape develop in the shear layer. They play a specific role with respect to local turbulent mixing in the ground surface boundary layer. Such curved slope constitutes a realistic model for alpine orography. We provide a novel procedure based on local turbulence anisotropy to track Görtler vortices for *in situ* measurements.

Ideal topography



$$G(x) = \frac{\langle u \rangle_{max}(x) \delta(x)}{\nu_{LES}(x)} \sqrt{\frac{\delta(x)}{R_{min}}}$$

Statistical Results



Prandtl Model

Analytical solution for the Prandtl model:

$$u_p(z) = V_o \sin(z/L_o) e^{-z/L_o}$$

$$\theta_p(z) - \theta_{ref}(z) = \Theta_o \cos(z/L_o) e^{-z/L_o}$$

Three characteristic scales L_o , V_o and Θ_o have to be prescribed from the boundary and ambient conditions. For heat flux boundary conditions, replacing viscous by turbulent quantities, and assuming mixing coefficients K_m and K_h constant along z , one gets

$$L_o = \frac{1}{Pr_t^{1/4}} \sqrt{\frac{2 K_m}{N_{ref} \sin(\alpha)}}$$

$$V_o = Pr_t^{1/4} \frac{\sqrt{2} g F_s}{\theta_o \sqrt{K_m N_{ref}^3 \sin(\alpha)}}$$

$$\Theta_o = Pr_t^{3/4} \frac{\sqrt{2} F_s}{\sqrt{K_m N_{ref} \sin(\alpha)}}$$

where $F_s = \frac{H_s}{\rho C_p}$ is the heat flux at the ground surface.

Present model adaptation with non constant diffusion:

$$K_m^{L_o} = K_m^{min} \left(1 + \frac{\partial K_m}{\partial z} \Big|_{L_o} z\right)$$

$$K_m^{V_o} = K_m^{min} \left(1 + \frac{\partial K_m}{\partial z} \Big|_{V_o} z\right)$$

$$K_m^{\Theta_o} = K_m^{min} \left(1 + \frac{\partial K_m}{\partial z} \Big|_{\Theta_o} z\right)$$

Prandtl model	K_m^{min}	$\frac{\partial K_m}{\partial z} \Big _{L_o}$	$\frac{\partial K_m}{\partial z} \Big _{V_o}$	$\frac{\partial K_m}{\partial z} \Big _{\Theta_o}$
$x = 1,450m$	$0.15m^2/s$	$0.04m/s$	$0.14m/s$	$0.17m/s$
$x = 1,500m$	$0.15m^2/s$	$0.05m/s$	$0.25m/s$	$0.33m/s$
$x = 1,550m$	$0.15m^2/s$	$0.06m/s$	$0.42m/s$	$0.58m/s$

CASE A1

Aknowledgements

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Lab OSUG@2020



Numerical Model

Meso-NH Model: CNRM & LA Toulouse France

- Pseudo-incompressible Navier-Stokes equations
- Anelastic approximation
- Buoyancy effects (gravity)
- No Coriolis effects
- Dry air (perfect gas)
- LES: e_{SGS} equation & mixing length closure
- Grid vertical refinement near the ground surface
- 5 \bar{M} grid points on 128 MPI proc. of IBM-SP6
- Initial conditions: air at rest with a constant $\frac{\partial \theta_{ref}}{\partial z} = N_{ref}^2 \theta_o$
- Ground surface cooling $H_s < 0$

Resolution/Boundary conditions

L_x (m)	L_y (m)	L_z (m)	n_x	n_y	n_z
3,200	1,280	7,250	128	128	300
Δx (m)	Δy (m)	Δz_{wall} (m)	Δz_{top} (m)	H (m)	α_{max}
25	10	1	120	1,000	35.5°

Case	u_*^{max} (m/s)	N_{ref} (s^{-1})	H_s (W/m^2)	$\langle u \rangle_{max}$ (m/s)
A0	0.19	0.011	-10	1.4
A1	0.24	0.013	-30	2.1
A2	0.18	0.013	-10	1.2
A3	0.15	0.020	-30	-
B1	0	0.013	-30	-

Anisotropy Invariant Map

Anisotropy tensor:

$$a_{ij} = \frac{\langle u_i u_j \rangle}{2 TKE} - \frac{1}{3} \delta_{ij}$$

second and third invariants:

$$I_2 = -\frac{1}{2} A^2 = -\frac{a_{ij} a_{ji}}{2}$$

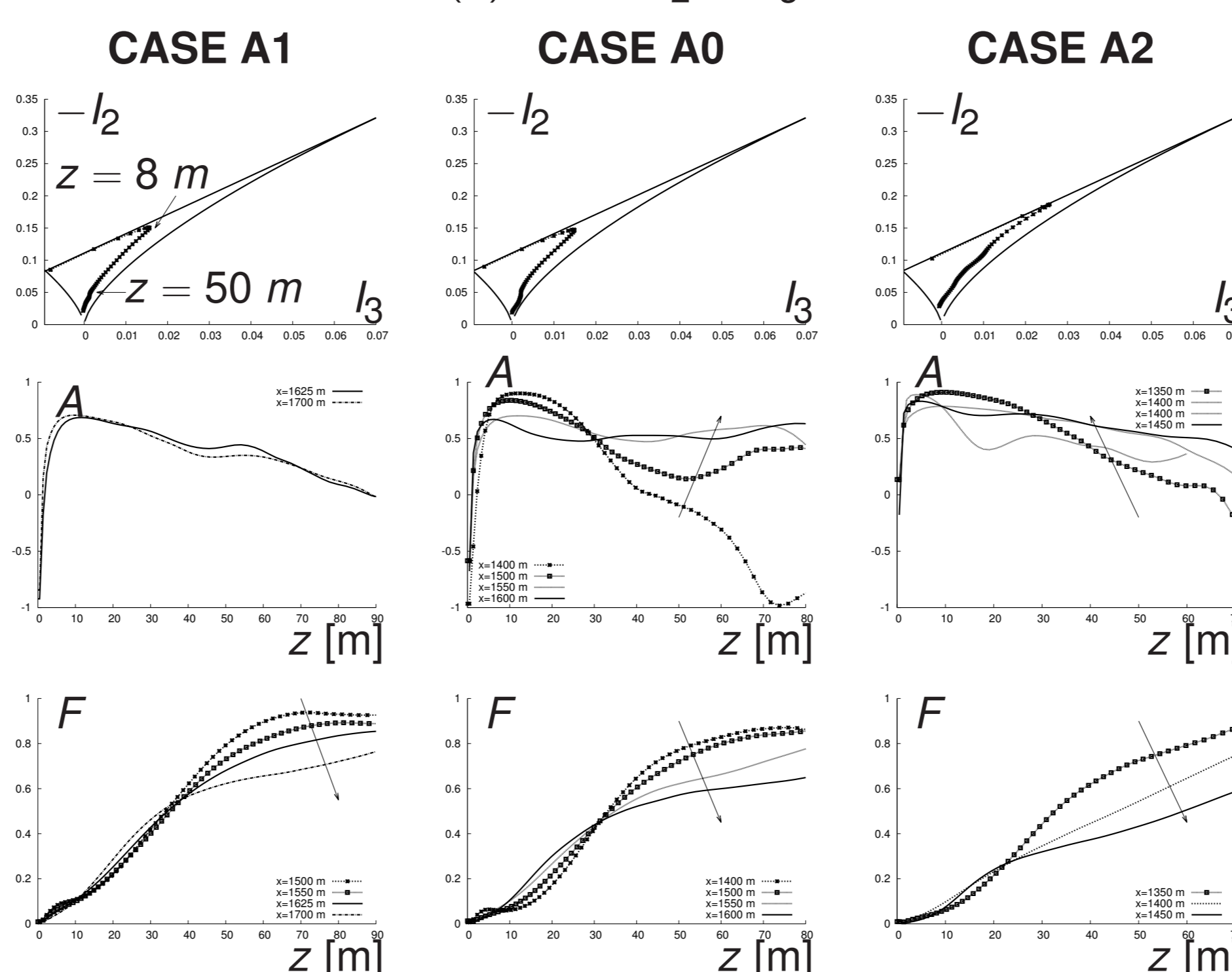
$$I_3 = \frac{1}{3} A^3 = \frac{a_{ij} a_{jk} a_{ki}}{3}$$

axisymmetric parameter:

$$A(z) = \frac{I_3}{2(-I_2/3)^{3/2}}$$

departure from isotropic turbulence:

$$F(z) = 1 + 9I_2 + 3I_3$$



TKE Budget

Turbulent Kinetic Energy

$$TKE = \frac{1}{2} (\langle u^2 \rangle + \langle v^2 \rangle + \langle w^2 \rangle)$$

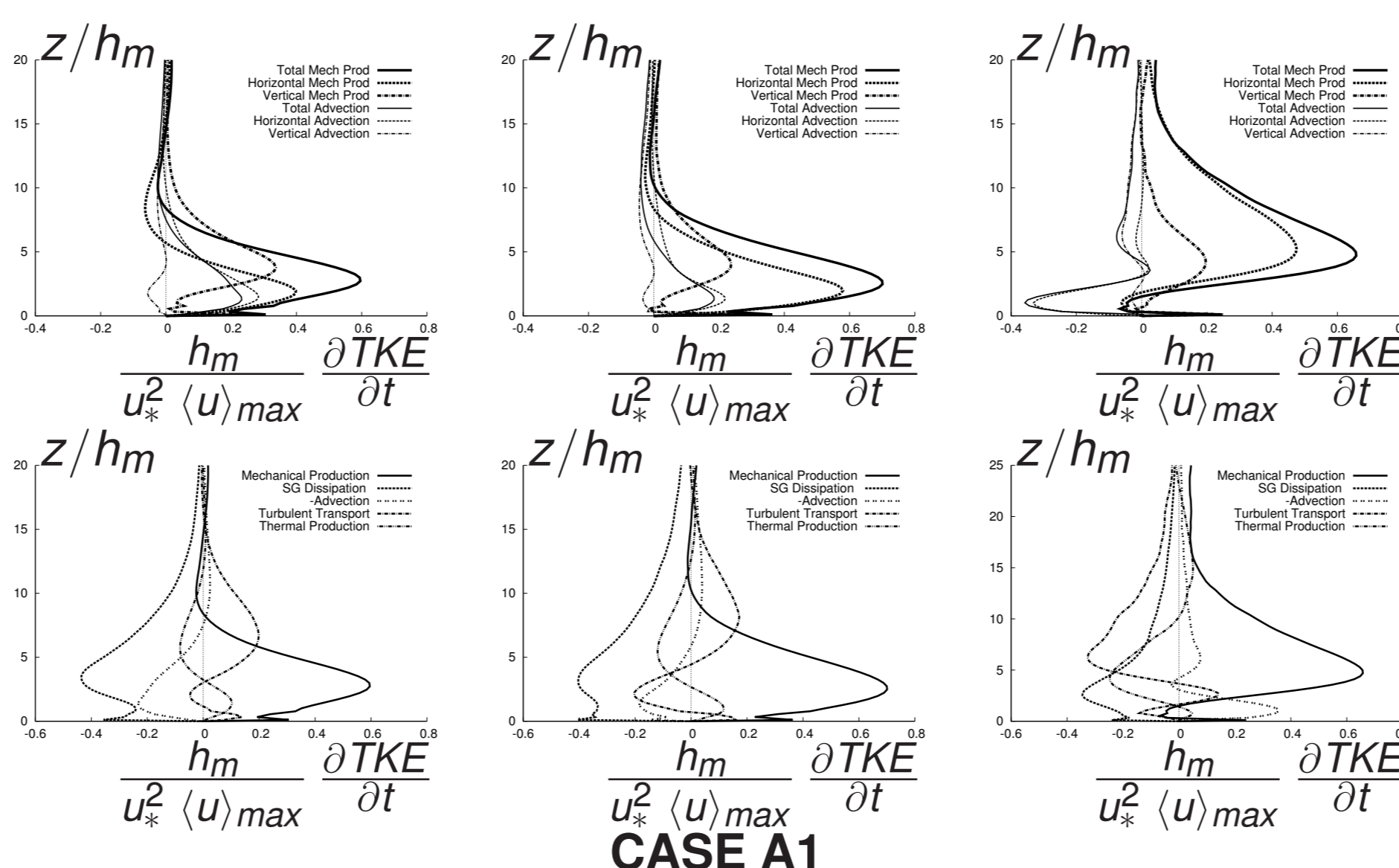
Horizontal and vertical production

$$\Pi_x = -\langle u' w' \rangle \frac{\partial \langle w \rangle}{\partial x} - \langle u'^2 \rangle \frac{\partial \langle u \rangle}{\partial x}$$

$$\Pi_z = -\langle u' w' \rangle \frac{\partial \langle u \rangle}{\partial z} - \langle w'^2 \rangle \frac{\partial \langle w \rangle}{\partial z}$$

horizontal and vertical Advection

$$Adv_x = \langle u \rangle \frac{\partial TKE}{\partial x} \quad Adv_z = \langle w \rangle \frac{\partial TKE}{\partial z}$$



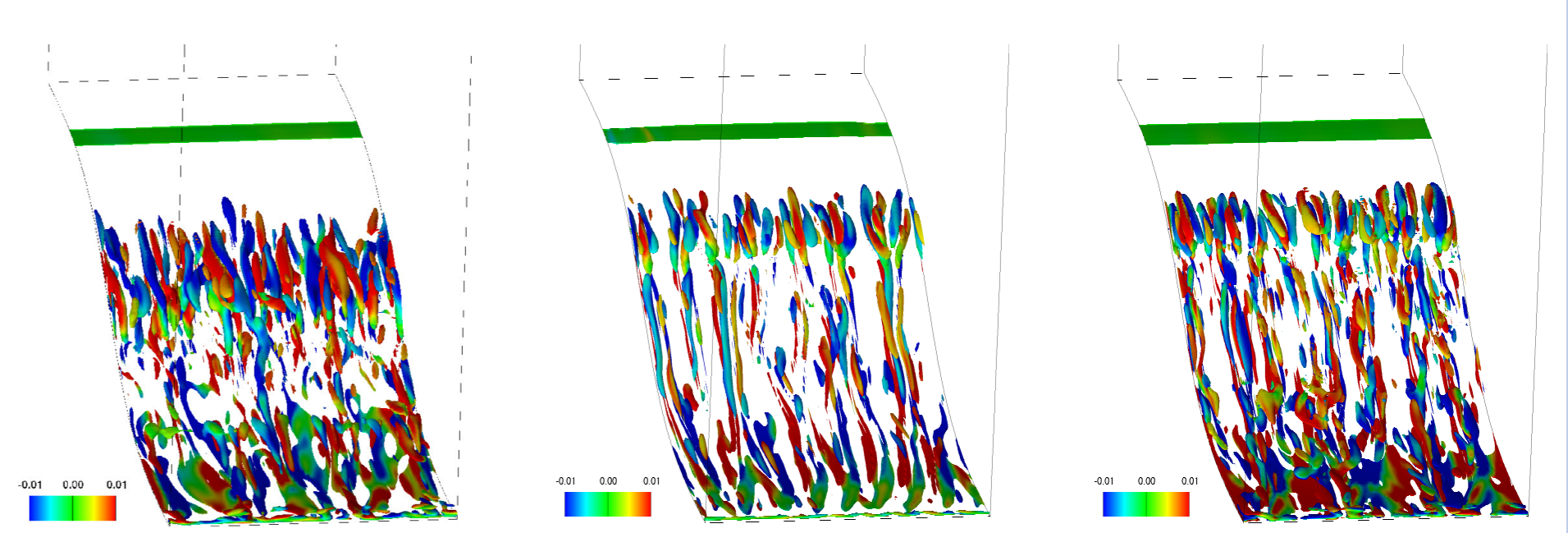
Communications

- Brun & Chollet Turbulent and Shear Flow Phenomena 2009
- Blein PhD Thesis 2016
- Brun, Blein & Chollet J. Atmospheric Sciences (submitted)

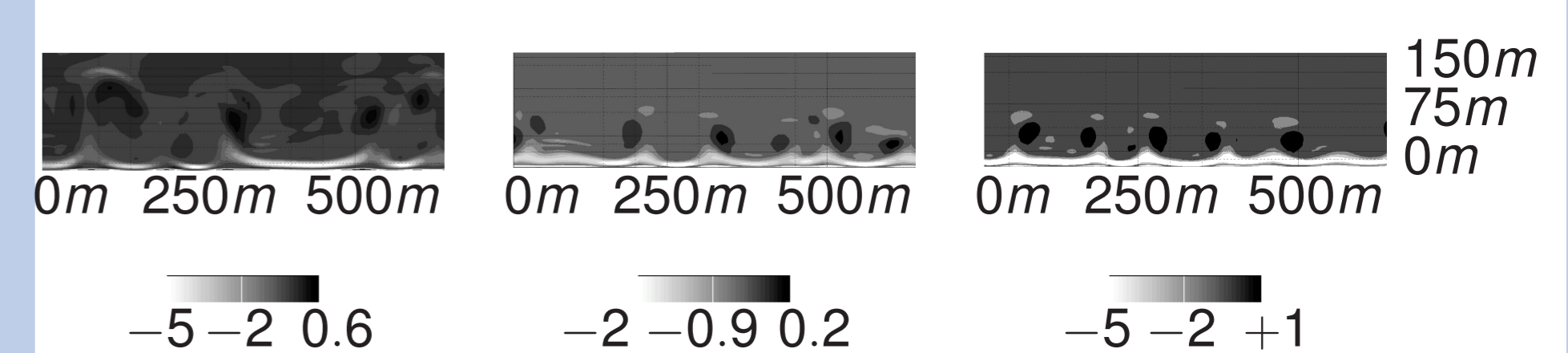
Flow Visualisation

- Transition to turbulence in the outer shear layer zone: $Ri \leq Ri_c = 0.2$
- Convex slope \rightarrow Görtler instability: spanwise wavelength $\lambda_y = 110 - 130 m$
- Strong streamwise counter-rotating (red/blue) vortices
- Local turbulent mixing enhanced

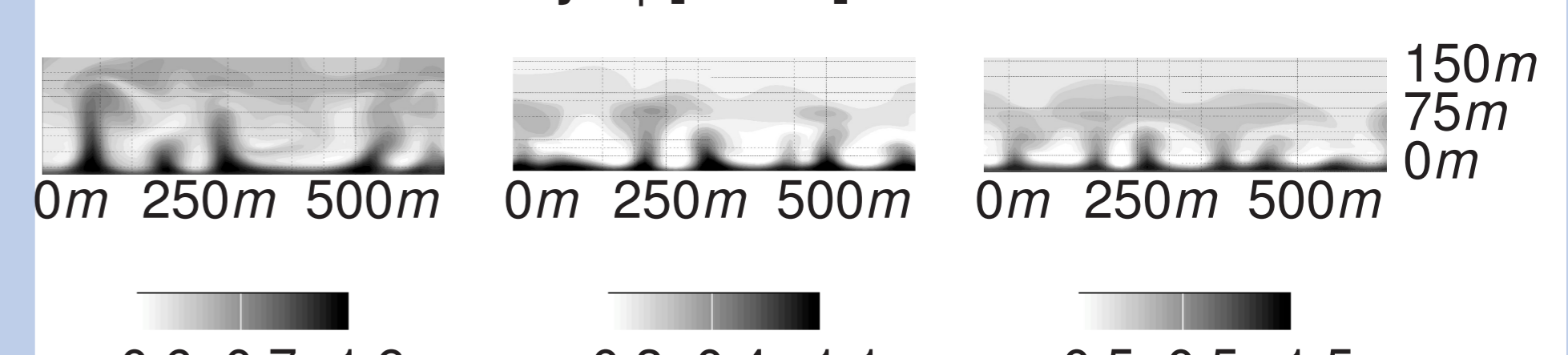
Q criterion $Q = \langle u \rangle_{max}^2 / H^2$



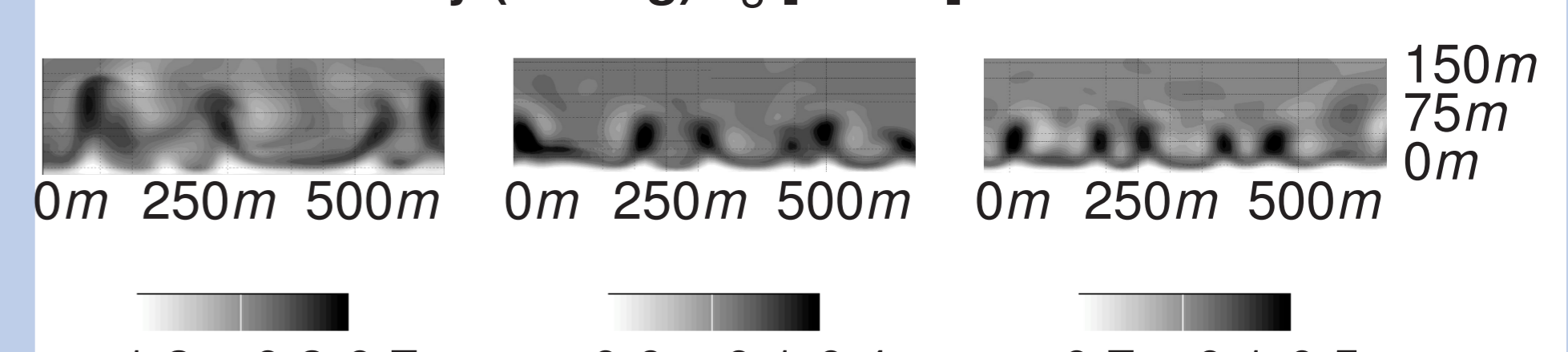
Q criterion $Q [10^{-3} s^{-2}]$



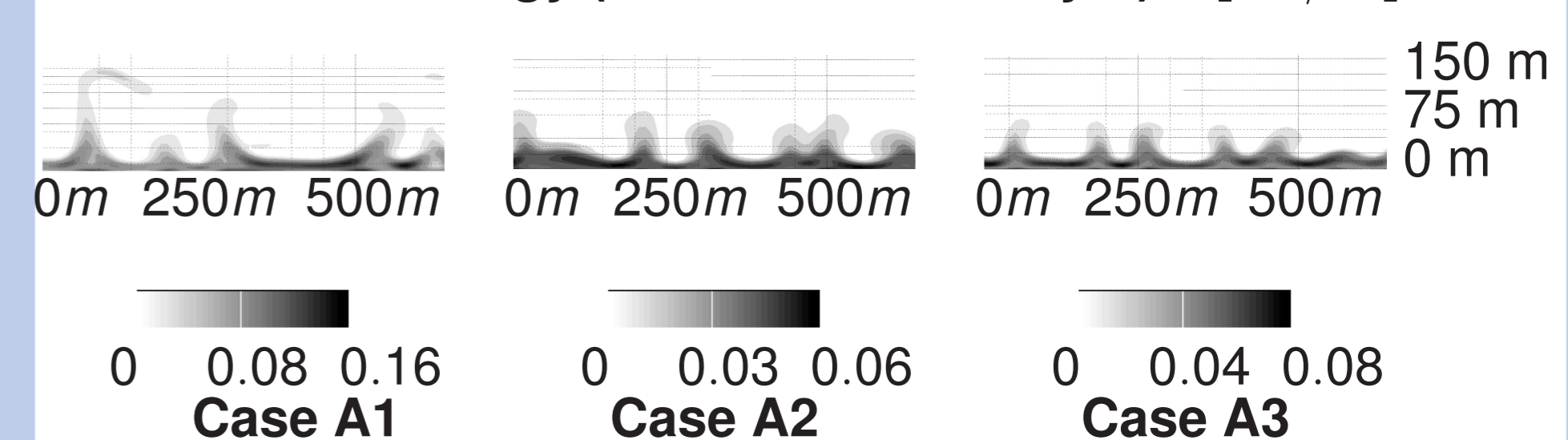
Streamwise velocity $u_1 [m.s^{-1}]$



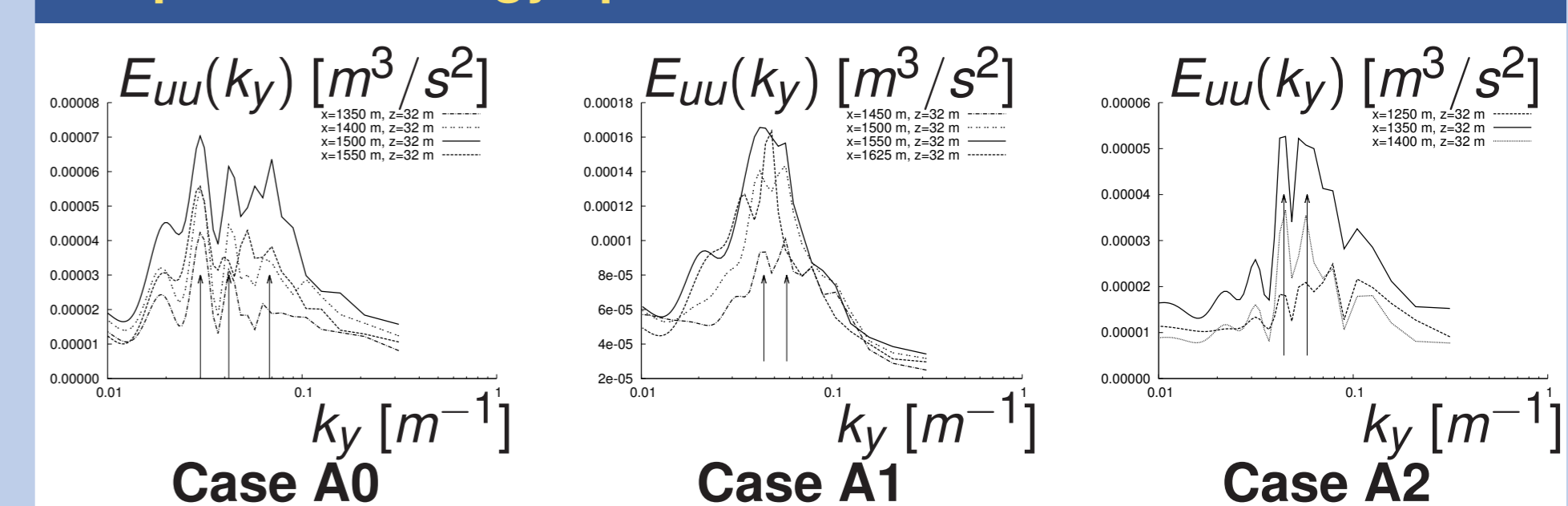
vertical velocity (mixing) $u_3 [m.s^{-1}]$



SGS Kinetic Energy (Turbulent shear layer) $e [m^2/s^2]$



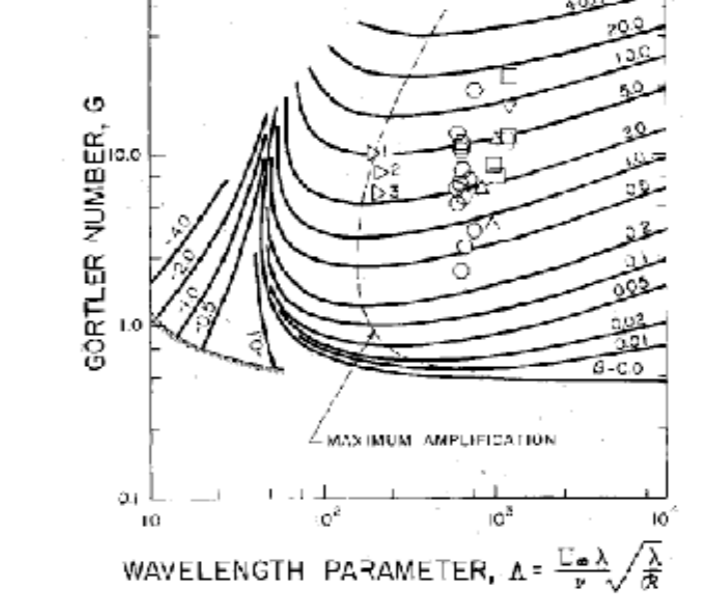
Spanwise Energy spectra



Görtler Stability Analysis

Boundary Layer along a concave wall

Katabatic jet along a convex surface



$$\Lambda(x) = \frac{\langle u \rangle_{max}(x) \lambda_y(x)}{\nu_{LES}(x)} \left(\frac{\lambda_y(x)}{R_{min}} \right)^{0.5}$$

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- Cuxart & Jimenez, Mixing processes in a nocturnal low-level jet: an LES study. *J. Atmospheric Sciences*, 2006
- Fedorovich & Shapiro, Structure of numerically simulated katabatic and anabatic flows along steep slopes. *Acta Geophysica*, 2009.
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