

ENERGETIC BUDGET OF DIRECT NUMERICAL SIMULATIONS IN A TURBULENT STRATIFIED FLOW.

Ernesto Horne^{1*}, Alexandre Delache², & Louis Gostiaux¹

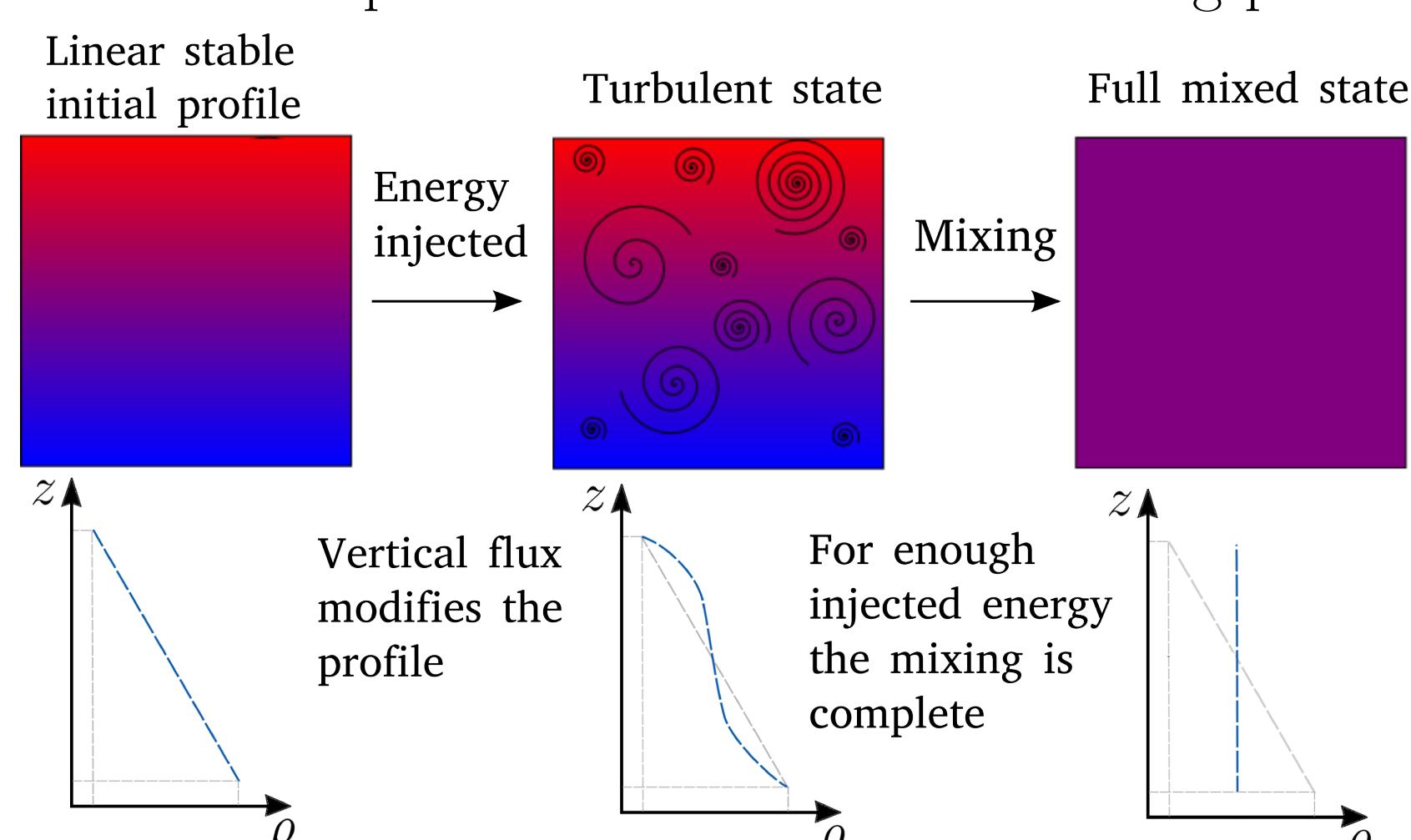
¹ LMFA UMR CNRS 5509, Ecole Centrale de Lyon, Université de Lyon, France

² Univ Lyon, UJM-Saint-Etienne, LMFA site de Saint Etienne UMR CNRS 5509, F-42023, Saint-Etienne, France

*ernesto.horne@ec-lyon.fr

Stably stratified environment

The vertical profile under a turbulent mixing process:



How does the vertical density profile evolves under turbulent mixing?

Global balance: advection and diffusion

For a stationary density profile the local vertical advection and diffusion must be globally conserved

$$w \partial_z \rho = \partial_z (K \partial_z \rho)$$

w : vertical upwelling

K : turbulent diffusivity

ρ : density

Osborn (1980) $\varepsilon_p = \eta \cdot \varepsilon = K \cdot N^2$ and $\eta = 1/6$

ε_p : fraction of the dissipation available to mix the fluid

$N^2 = -\frac{g}{\rho} \partial_z \rho$: buoyancy frequency

η : mixing efficiency

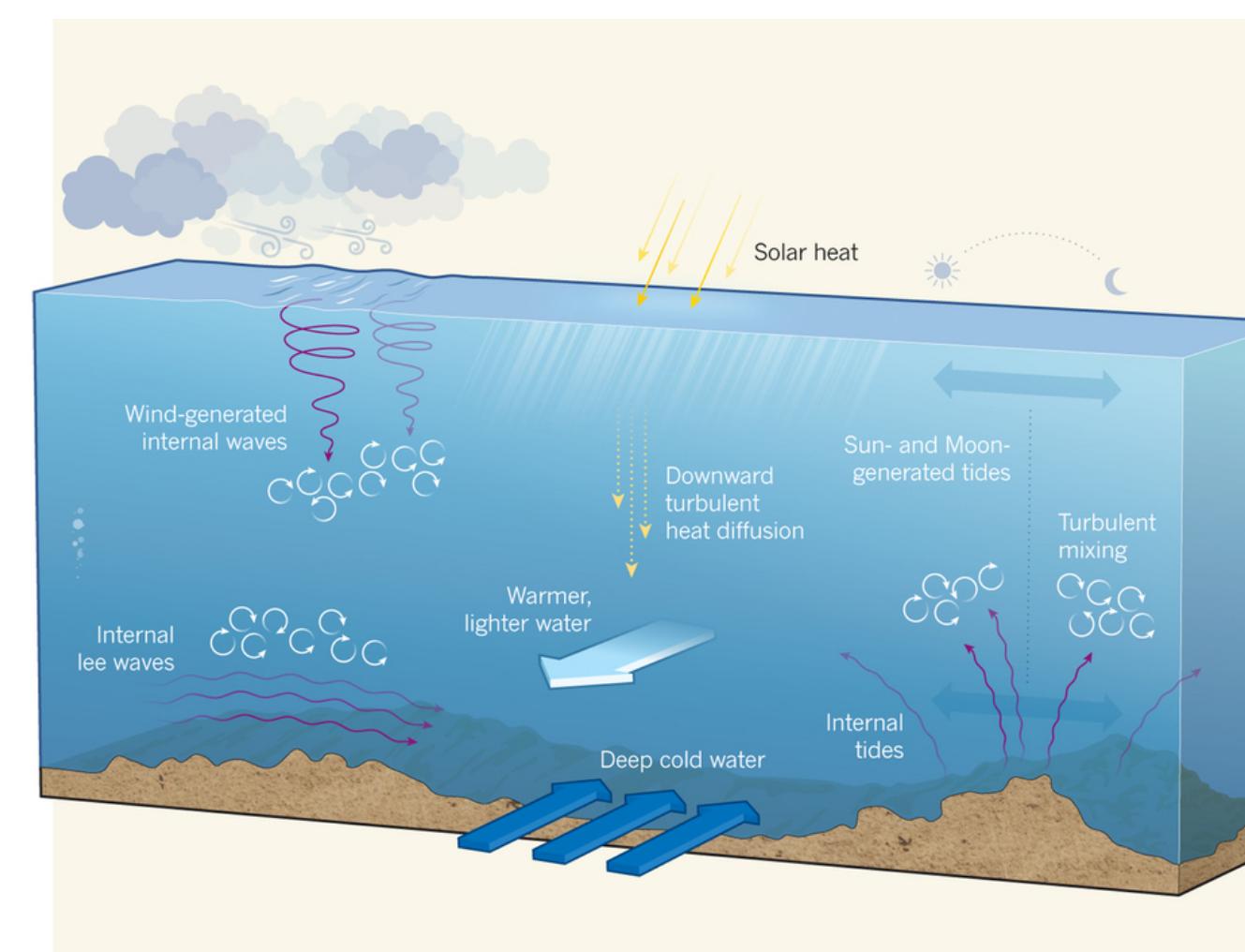
ε : viscous dissipation

In the limit of a homogeneous fluid, mixing cannot drive a buoyancy flux

$$K \cdot N^2 = \eta \cdot \varepsilon \rightarrow 0$$

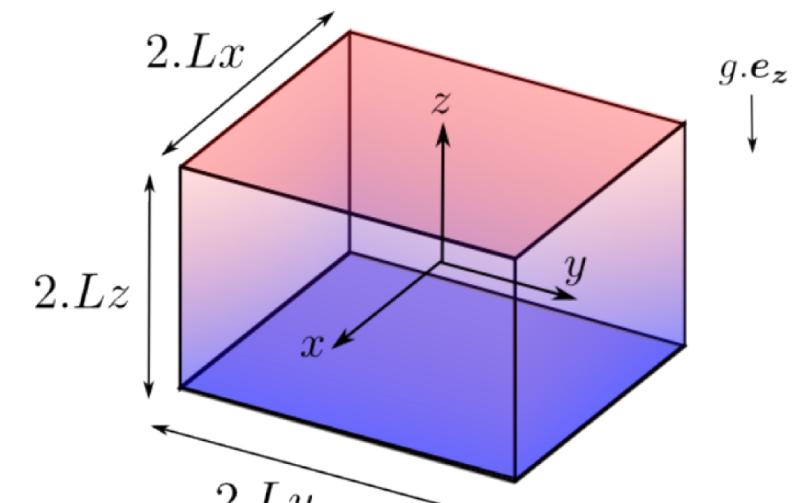
MacKinnon 2013. Nature.

Motivation: the Ocean



- Many different processes generate mixing in the ocean interior.
- Mixing drives downward transport of heat and dissolved greenhouse gases, and upward transport of biological nutrients.
- Drives a global network of currents (MOC).
- Induces meridional heat transport.
- Important effect on climate.

Method: Direct Numerical Simulations



- Navier-Stokes under Boussinesq approximation.
- Standard pseudospectral method (Fourier).
- Velocity field imposed for initial conditions.
- Penalisation region in the boundaries.
- Background density can evolve.

Equations: Dimensionless Navier-Stokes

$$\begin{aligned} \frac{\partial \mathbf{u}}{\partial t} - \omega \times \mathbf{u} &= -\nabla p - \theta \cdot \mathbf{Ri} \cdot z + \frac{1}{Re} \cdot \nabla^2 \mathbf{u} \\ \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \theta &= \frac{1}{Re \cdot Sc} \nabla^2 \theta \\ \nabla \cdot \mathbf{u} &= 0 \end{aligned}$$

Dimensionless numbers:

$$Ri = \frac{\theta_0 \cdot g \cdot L_0}{U_0^2}, Re = \frac{L_0 \cdot U_0}{\nu}, Sc = \frac{\nu}{\kappa}$$

Turbulent integral scales

$$L_0, U_0$$

DNS parameters:

$$N_{grid} = 1024, Re = 1000, Sc = 5$$

run	1	2	3	4	5	6	7
Ri	2	10	44	180	740	3000	11000

Reduced buoyancy

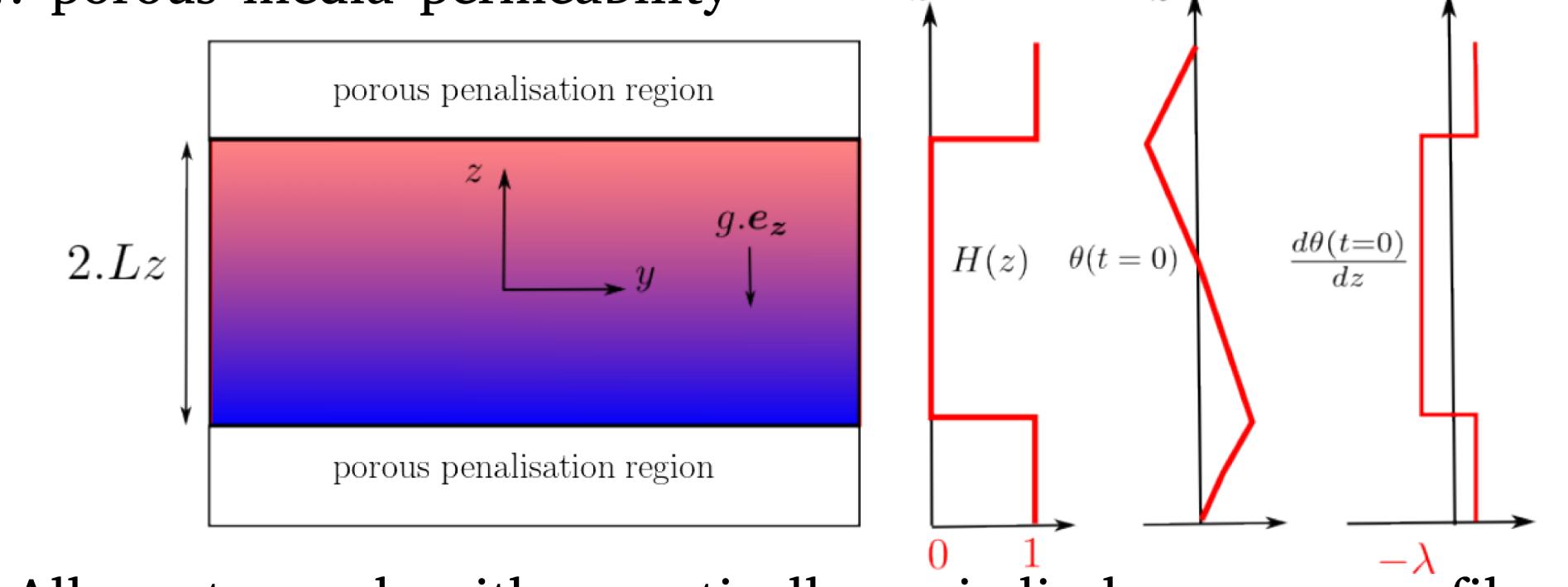
$$\theta = \frac{\rho - \rho_0}{\rho_0}$$

$$\theta_0 = \frac{\rho^{max} - \rho^{min}}{\rho_0}$$

In the penalisation region: Inspired by Kadosh et al. 2012

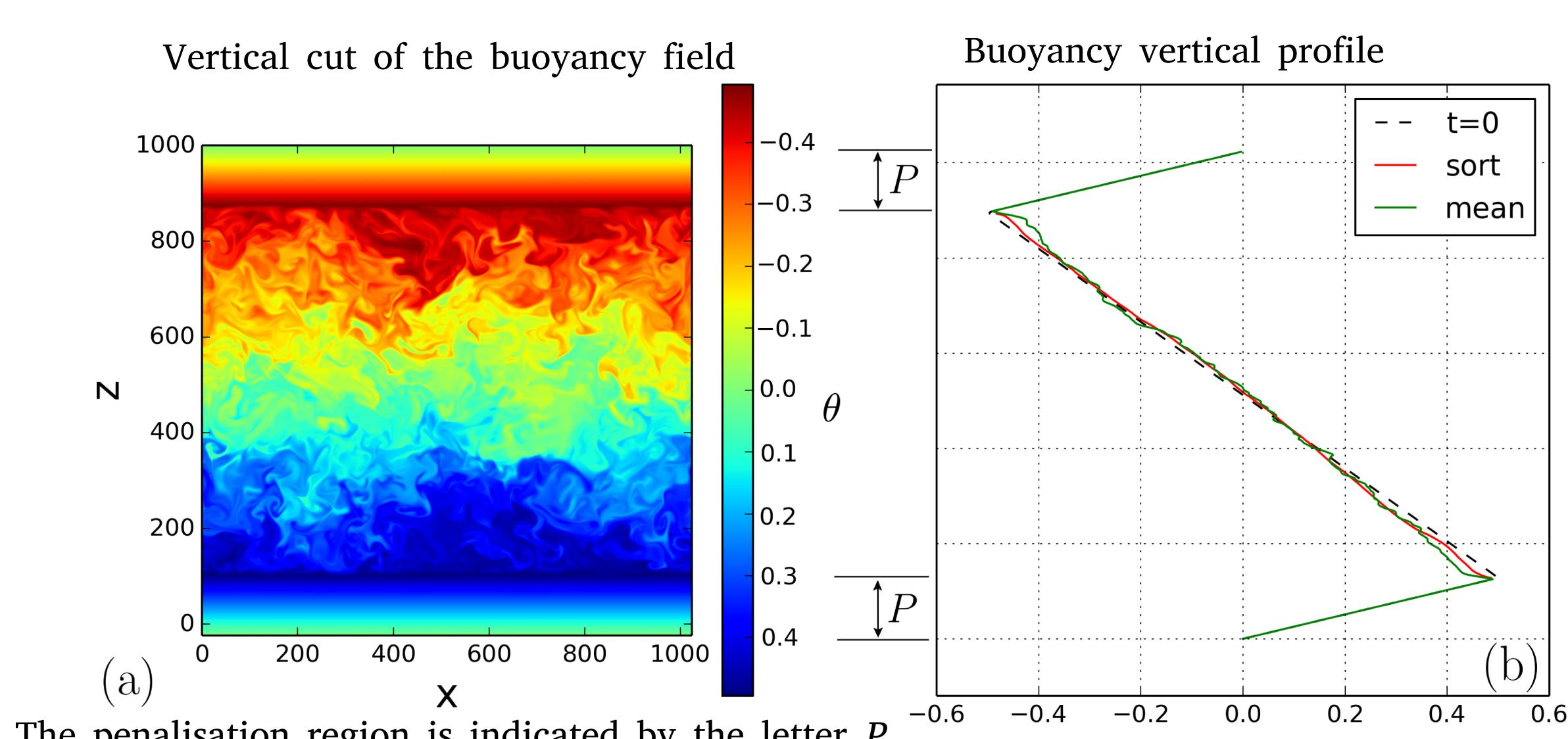
$$\begin{aligned} \frac{\partial \mathbf{u}}{\partial t} - \omega \times \mathbf{u} &= \nabla p - \theta \cdot \mathbf{Ri} \cdot (z) \cdot (1 - H) + \frac{1}{Re} \cdot \nabla^2 - \frac{1}{\eta_p} \cdot H \cdot \mathbf{u} \\ \frac{\partial \theta}{\partial t} &= \nabla \cdot [(1 - H) \cdot \mathbf{u} \cdot \theta] + \nabla \cdot \left\{ \left[\frac{1}{Re \cdot Sc} (1 - H) + \eta_p \cdot H \right] \cdot \nabla \theta \right\} \end{aligned}$$

η_p : porous media permeability



Allows to work with a vertically periodic buoyancy profile

Results: buoyancy field



The penalisation region is indicated by the letter P

Buoyancy fields:

- Horizontal average buoyancy field $\langle \theta(x, y, z, t) \rangle_{x,y}$
- Horizontal average sorted buoyancy field $\langle \theta_S(x, y, z, t) \rangle_{x,y}$

The sorted buoyancy field is obtained by sorting every element of fluid of the buoyancy field. It allows to measure instantaneously the irreversible mixing.

- Initially, the buoyancy field is sorted $\theta(x, y, z, t = 0) = \theta_S(x, y, z, t = 0)$

Potential energies

Potential energy:

$$E_p = g \cdot \rho_0 \int \theta \cdot z \cdot dV$$

Given by the instantaneous center of mass of the fluid

Background potential energy:

$$E_b = g \cdot \rho_0 \int \theta_S \cdot z \cdot dV$$

Indicates the irreversible mixing done in the fluid

Available potential energy:

$$E_{ape} = g \cdot \rho_0 \int (\theta - \theta_S) \cdot z \cdot dV$$

Represents the instantaneous potential energy that has not yet been transformed in irreversible mixing

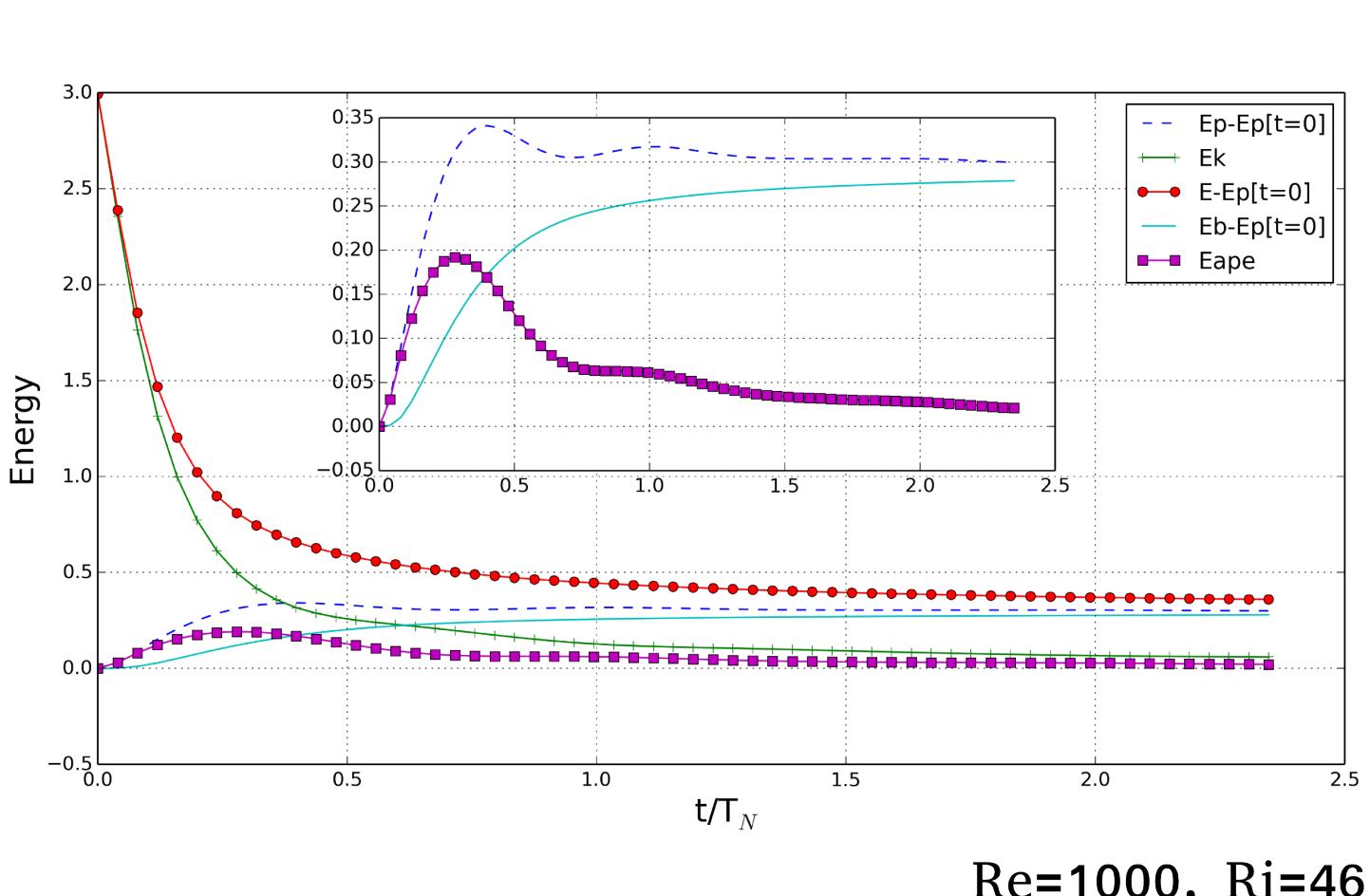
Kinetic energy:

$$E_k = \frac{\rho_0}{2} \int \mathbf{u}^2 \cdot dV$$

Winters 1995:

Energies: Evolution

Evolution of potential and kinetic energies:

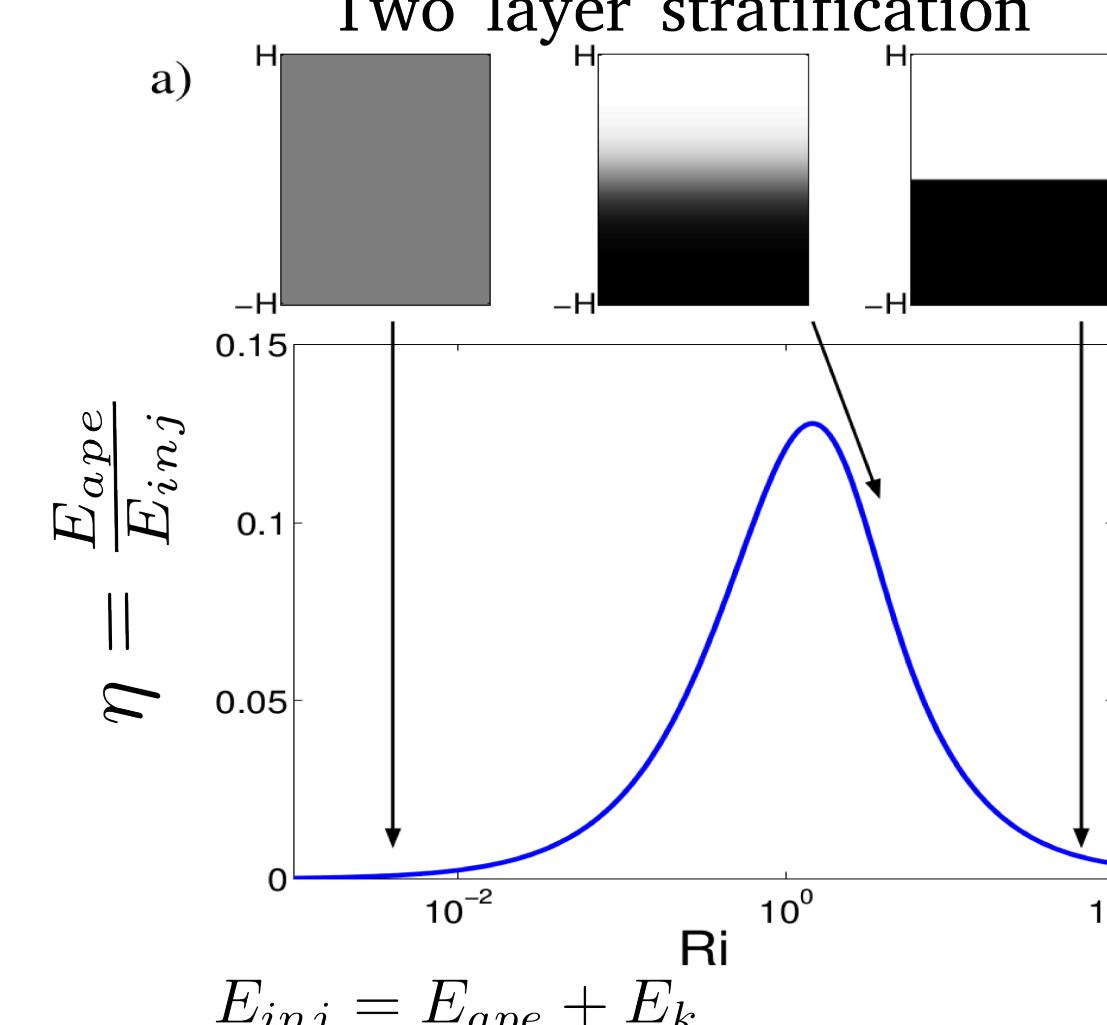


Three phases of the flow:

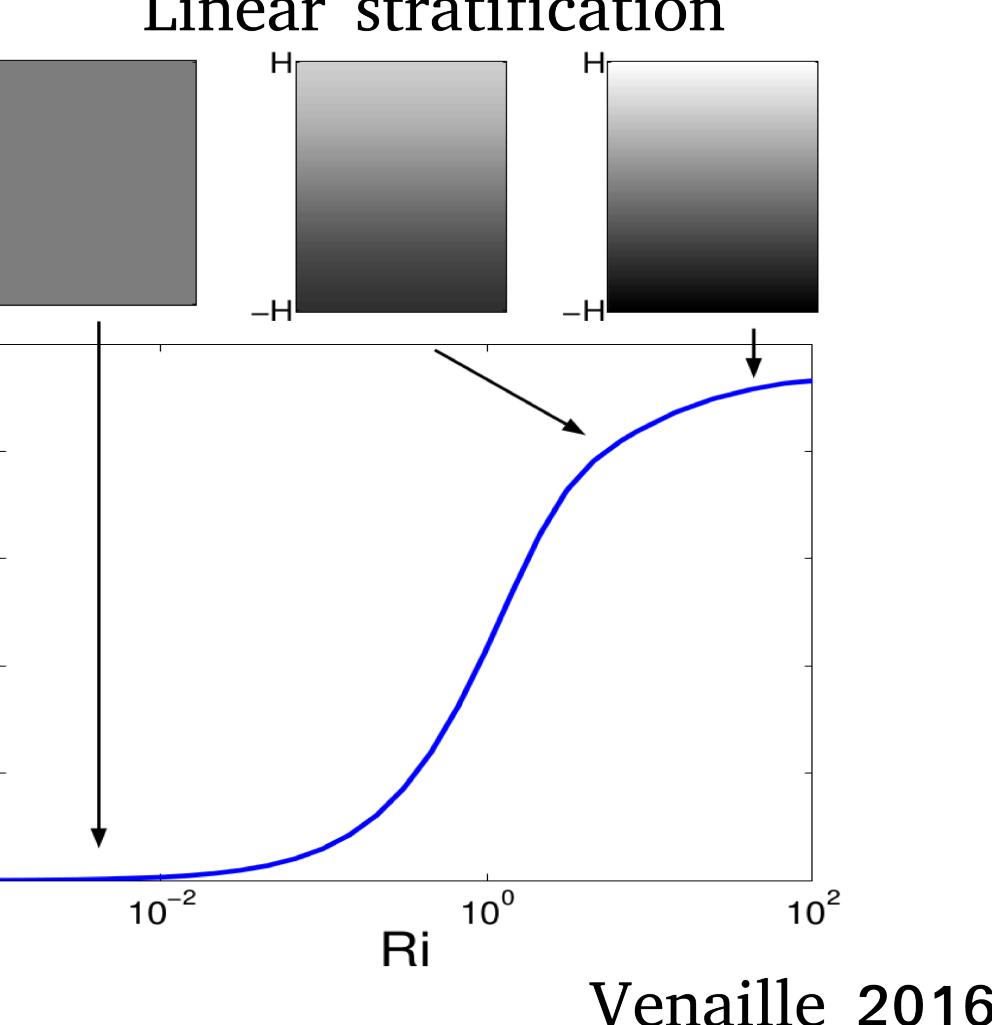
- Initially the energy is injected through kinetic energy
- Once the turbulence is developed, the kinetic energy decays and, therefore the rate of irreversible mixing decreases.
- Finally, turbulence is weaker and the flow transitions towards laminarization.

Perspectives: mixing efficiency through a statistical theory

Two layer stratification



Linear stratification



Venaille 2016.

Bibliography

- MacKinnon, J.: Oceanography: Mountain waves in the deep ocean. *Nature* 501, 321-322 (2013).
 Kadosh, B. et al.: A volume penalization method for incompressible flows and scalar advection-diffusion with moving obstacles. *JCP* 231, 4365-4383 (2012).
 Winters, K. B. et al.: Available potential energy and mixing in density-stratified fluids. *JFM* 289, 115-128 (1995).
 Venaille, A. et al.: A statistical theory of mixing in stratified fluids. Submitted to *JFM* (2016).